

## Appendix B

### Analysis of Scattered Waves on Three Layers of Tree Trunk

(TM mode)

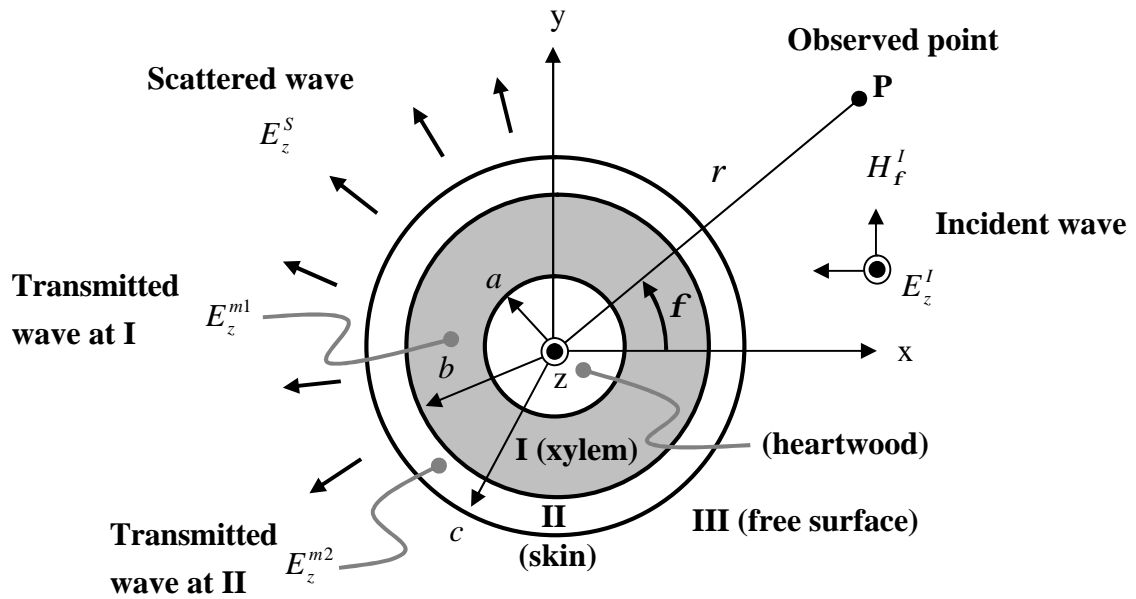


Figure B.1. Geometry of analysis of scattered TM mode wave (three layers)

The two-dimensional model of tree trunk is shown in figure B.1. Three layers of media compose this model of tree trunk with infinite length in z-axis. The radii of heartwood, xylem and skin layer are  $a$ ,  $b$ , and  $c$ , respectively. The properties of xylem and skin are determined by  $\mathbf{e}_{r1}$ ,  $\mathbf{m}_{r1}$  and  $\mathbf{e}_{r2}$ ,  $\mathbf{m}_{r2}$ , respectively. Where,  $\mathbf{e}_{ri}$  and  $\mathbf{m}_{ri}$  ( $i=1,2$ ) are complex dielectric constant and complex permeability, respectively. The water content of heartwood is high, consequently, heartwood may be assumed to be an infinite length of

perfect conductor or electromagnetic fields in heartwood is zero. Here, incident wave is assumed as a plane wave that has transverse magnetic (TM) mode and incident angle  $\mathbf{f}$  with respect to direction of observed point P from origin of coordinate. This wave propagates in  $-x$  direction. Based on this figure, the  $z$  component of the electromagnetic fields in free space, xylem and skin are determined as

$$\text{Incident wave} \quad E_z^I = E_o^I \exp(jk_o r) = E_o^I \sum_{m=0}^{\infty} U_m J_m(k_o r) j^m \cos m\mathbf{f} \quad (r > c) \quad (\text{B.1})$$

$$\text{Scattered wave} \quad E_z^S = E_o^I \sum_{m=0}^{\infty} b_m H_m^{(2)}(k_o r) \cos m\mathbf{f} \quad (r > c) \quad (\text{B.2})$$

$$\text{Medium I} \quad E_z^{m1} = E_o^I \sum_{m=0}^{\infty} \{a_{1m} J_m(k_1 r) + a'_{1m} N_m(k_1 r)\} \cos m\mathbf{f} \quad (a < r \leq b) \quad (\text{B.3})$$

$$\text{Medium II} \quad E_z^{m2} = E_o^I \sum_{m=0}^{\infty} \{a_{2m} J_m(k_2 r) + a'_{2m} N_m(k_2 r)\} \cos m\mathbf{f} \quad (b < r \leq c) \quad (\text{B.4})$$

Where the wave numbers of each medium are  $k_1 = k_o \sqrt{\mathbf{m}_{r1} \mathbf{e}_{r1}}$  and  $k_2 = k_o \sqrt{\mathbf{m}_{r2} \mathbf{e}_{r2}}$ , and  $k_o$  is wave number in free space.  $a_{1m}$  to  $b_m$  are amplitude coefficients.  $E_o^I$  is amplitude coefficient of incident electric field.  $J_m$ ,  $N_m$ , and  $H_m^{(2)}$  are  $m$ -th of Bessel function, Neumann function, and 2nd kind of Hankel function, respectively. Then  $m = 0, 1, 2, 3, \dots$  and

$$U_m = \begin{cases} 1, & m = 0 \\ 2, & m = 1, 2, \dots \end{cases} \quad (\text{B.5})$$

By substituting (B.1) to (B.4) in Maxwell's equations below

$$\nabla \times \mathbf{E} = -\mathbf{m} \frac{\partial \mathbf{H}}{\partial t} \quad (\text{B.6})$$

the magnetic field of each medium is derived as

$$\text{Incident wave} \quad H_f^I = \frac{k_o E_o^I}{j\omega \mathbf{m}_o} \sum_{m=0}^{\infty} U_m J'_m(k_o r) j^m \cos m\mathbf{f} \quad (r > b) \quad (\text{B.7})$$

$$\text{Scattered wave } H_f^S = \frac{k_o E_o^I}{j\omega \mathbf{m}_o} \sum_{m=0}^{\infty} b_m H_m^{(2)'}(k_o r) \cos m\mathbf{f} \quad (r > b) \quad (\text{B.8})$$

$$\text{Medium I } H_f^{m1} = \frac{k_1 E_o^I}{j\omega \mathbf{m}_1} \sum_{m=0}^{\infty} \{a_{1m} J_m'(k_1 r) + a'_{1m} N_m'(k_1 r)\} \cos m\mathbf{f} \quad (c < r \leq a) \quad (\text{B.9})$$

$$\text{Medium II } H_f^{m2} = \frac{k_2 E_o^I}{j\omega \mathbf{m}_2} \sum_{m=0}^{\infty} \{a_{2m} J_m'(k_2 r) + a'_{2m} N_m'(k_2 r)\} \cos m\mathbf{f} \quad (a < r \leq b) \quad (\text{B.10})$$

Further, by substituting (B.1) to (B.4) and (B.7) to (B.10) into the boundary condition of each interface between media given below:

$$r = a \quad E_z^{m1} = 0 \quad (\text{B.11})$$

$$r = b \quad H_f^{m1} = H_f^{m2} \quad (\text{B.12})$$

$$E_z^{m1} = E_z^{m2} \quad (\text{B.13})$$

$$r = c \quad H_f^{m2} = H_f^S + H_f^I \quad (\text{B.14})$$

$$E_z^{m2} = E_z^S + E_z^I \quad (\text{B.15})$$

the amplitude coefficient  $b_m$  of scattered wave from tree trunk  $E_f^S$  is obtained as;

$$b_m = \frac{U_m j^m \{a_{6m} J_m'(k_o c) - a_{7m} J_m(k_o c)\}}{a_{7m} H_m^{(2)}(k_o c) - a_{6m} H_m^{(2)'}(k_o c)} \quad (\text{B.16})$$

where

$$a_{7m} = \frac{k_2 \mathbf{m}_o}{k_o \mathbf{m}_2} \left\{ N_m'(k_2 c) - \frac{a_{5m}}{a_{4m}} J_m'(k_2 c) \right\} \quad (\text{B.17})$$

$$a_{6m} = N_m(k_2 c) - \frac{a_{5m}}{a_{4m}} J_m(k_2 c) \quad (\text{B.18})$$

$$a_{5m} = a_{3m} \frac{k_1}{\mathbf{m}_1} N_m(k_2 b) - a_{2m} \frac{k_2}{\mathbf{m}_2} N_m'(k_2 b) \quad (\text{B.19})$$

$$a_{4m} = a_{3m} \frac{k_1}{\mathbf{m}_1} J_m(k_2 b) - a_{2m} \frac{k_2}{\mathbf{m}_2} J_m'(k_2 b) \quad (\text{B.20})$$

$$a_{3m} = N_m'(k_1 b) - a_{1m} J_m'(k_1 b) \quad (\text{B.21})$$

$$\mathbf{a}_{2m} = N_m(k_1 b) - \mathbf{a}_{1m} J_m(k_1 b) \quad (\text{B.22})$$

$$\mathbf{a}_{1m} = \frac{N_m(k_1 a)}{J_m(k_1 a)} \quad (\text{B.23})$$

Finally, by substituting the amplitude coefficient  $b_m$  of (B.16) into (B.2), the scattered electric field is obtained.